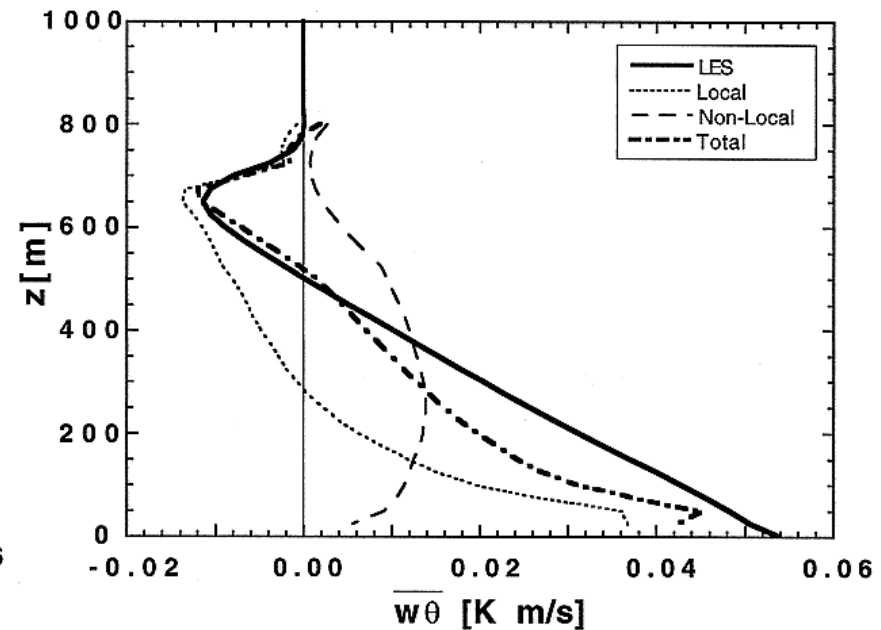
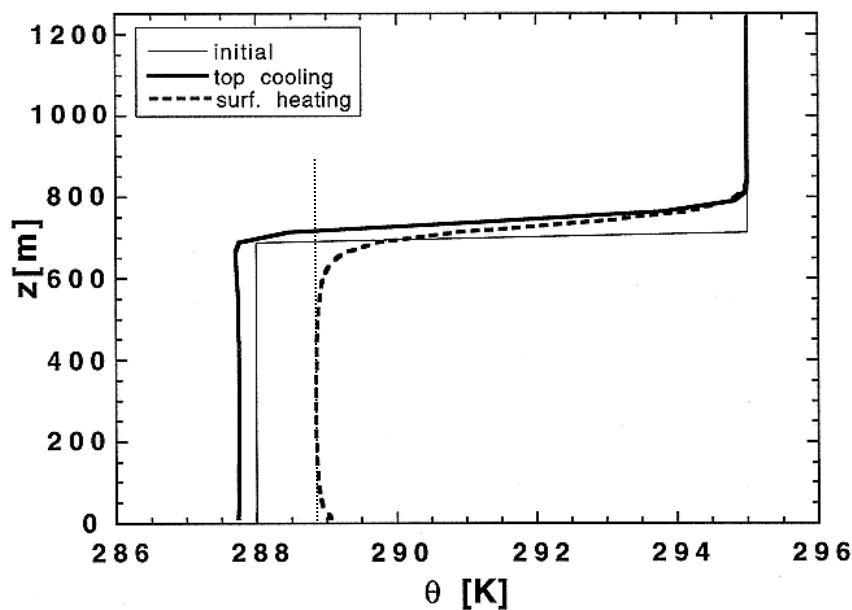


A challenge to downgradient diffusion: Countergradient heat transport

- In dry convective boundary layer, deep eddies transport heat
- This breaks correlation between local gradient and heat flux
- LES shows slight q min at $z=0.4h$, but $w'q' > 0$ at $z < 0.8h$
- ‘Countergradient’ heat flux for $0.4 < z/h < 0.8$...first recognized in 1960s by Telford, Deardorff, etc.



Cuijpers and Holtslag 1998

Nonlocal K-profile schemes

$$\overline{w'a'} = -K_a(z) \frac{\partial a}{\partial z} + \text{another 'nonlocal' term}$$

(Holtslag-Boville in CAM3/4, YSU in WRF, EDMF in ECMWF):

Derivation of nonlocal schemes

Heat flux budget:
$$\underbrace{\frac{\partial \overline{w'\theta'}}{\partial t}}_S = -\underbrace{\overline{w'w'}}_M \frac{\partial \bar{\theta}}{\partial z} - \underbrace{\frac{\partial \overline{w'w'\theta'}}{\partial z}}_T + \underbrace{\frac{g}{\theta_0} \overline{\theta'\theta'}}_B - \underbrace{\frac{1}{\rho_0} \overline{\theta' \frac{dp'}{dz}}}_P$$

Neglect storage S

Empirically:

$$T \approx B + 2 \frac{w_*^2 \theta_*}{h}$$

$$P = -aB - \frac{\overline{w'\theta'}}{\tau}$$

For convection, $a=0.5$, so

$$\overline{w'\theta'} = -\underbrace{\frac{\tau}{2} \overline{w'w'}}_{K_H(z)} \frac{\partial \bar{\theta}}{\partial z} + \tau \frac{w_*^2 \theta_*}{h}$$

Take $\tau = 0.5h/w_*$ to get zero θ gradient at $0.4h$.

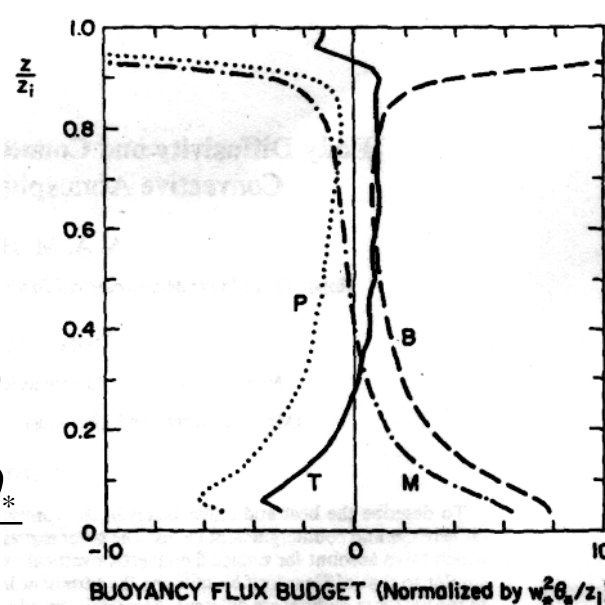


FIG. 1. The normalized terms at the rhs of the heat-flux equation (1), as a function of relative height (adopted from Moeng and Wyngaard 1989). The terms are defined in the text of section 2a.

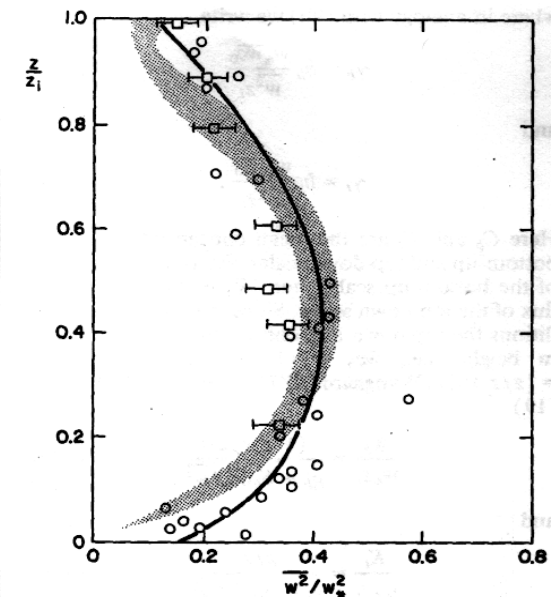


FIG. 4. The nondimensional vertical-velocity variance of (15a) (solid curve) in comparison with the (96)³ LES data (shaded area; Moeng and Wyngaard 1989), the AMTEX data (circles; Lenschow et al. 1980), and convection tank experiments (squares; Deardorff and Willis 1985).

Holtslag and Moeng (1991)

Nonlocal parameterization, continued

This has the form $\overline{w'\theta'} = -K_H(z) \left(\frac{\partial\theta}{\partial z} - \gamma_\theta \right)$ where $\gamma_\theta = \frac{2w_*^2\theta_*}{\overline{w'w'h}}$

Although the derivation suggests γ_θ is a strong function of z , the parameterization treats it as a constant evaluated at $z = 0.4h$ to obtain the correct heat flux there with $d\theta/dz = 0$:

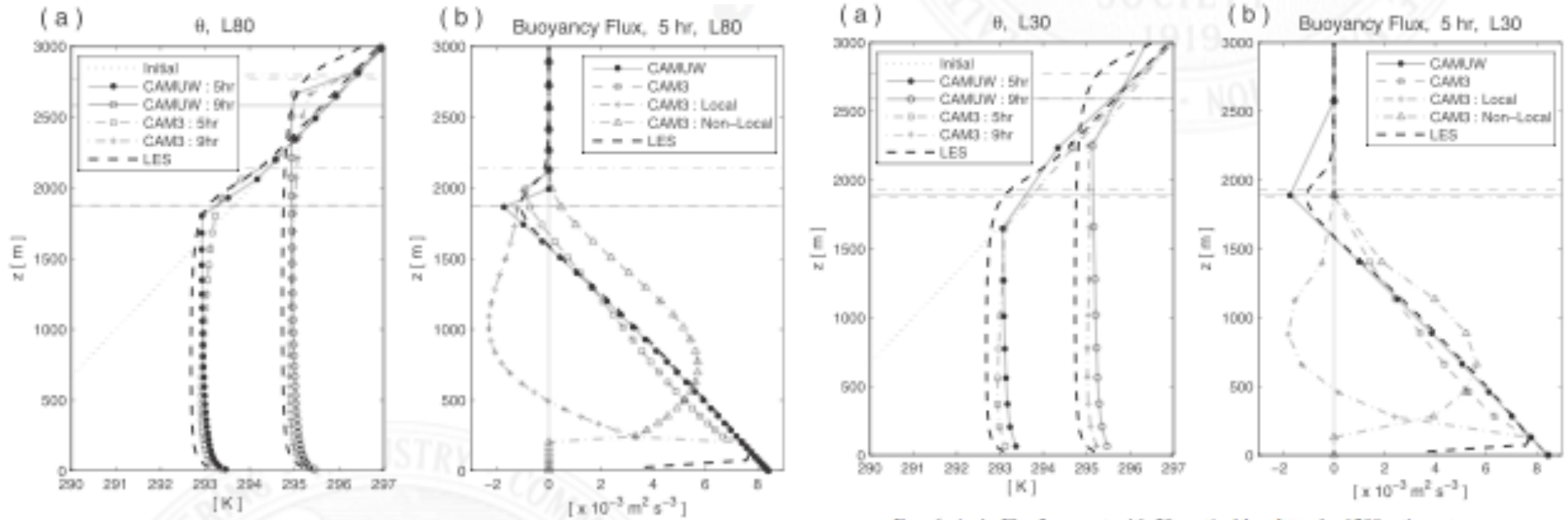
$$\overline{w'w'}(0.4h) = 0.4w_*^2 \Rightarrow \gamma_\theta = 5\theta_*/h.$$

The eddy diffusivity can be parameterized from vert. vel. var.:

$$\overline{w'w'}(z) = 2.8w_*^2Z(1-Z)^2, \quad Z = z/h \Rightarrow K_H(z) = 0.7w_*z(1-Z)^2$$

With cleverly chosen velocity scales, this can be seamlessly combined with a K-profile for stable BLs to give a generally applicable parameterization (Holtslag and Boville 1993).

CBL comparison



Bretherton and Park 2009

FIG. 6. As in Fig. 5, except with 30 vertical levels and a 1200-s time step.

- Sfc heating of 300 W m^{-2}
- No moisture or mean wind
- UW TKE scheme with entrainment closure and HB scheme give similar results at both high and low res.
- Overall, can get comparably good results from TKE and profile-based schemes on these archetypical cases.

EDMF

$$\overline{w'\phi'} \cong -K \frac{\partial \bar{\phi}}{\partial z}$$

$$\overline{w'\phi'} \cong M(\phi_u - \bar{\phi})$$

