

2 **Numerical modeling of the impact of sea-level rise on fringing**
3 **coral reef hydrodynamics and sediment transport**

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8 **Abstract** Most climate projections suggest that sea level
9 may rise on the order of 0.5–1.0 m by 2100; it is not clear,
10 however, how fluid flow and sediment dynamics on
11 exposed fringing reefs might change in response to this
12 rapid sea-level rise. Coupled hydrodynamic and sediment-
13 transport numerical modeling is consistent with recent
14 published results that suggest that an increase in water
15 depth on the order of 0.5–1.0 m on a 1–2 m deep exposed
16 fringing reef flat would result in larger significant wave
17 heights and setup, further elevating water depths on the
18 reef flat. Larger waves would generate higher near-bed
19 shear stresses, which, in turn, would result in an increase in
20 both the size and the quantity of sediment that can be
21 resuspended from the seabed or eroded from adjacent
22 coastal plain deposits. Greater wave- and wind-driven
23 currents would develop with increasing water depth,
24 increasing the alongshore and offshore flux of water and
25 sediment from the inner reef flat to the outer reef flat and
26 fore reef where coral growth is typically greatest. Sediment
27 residence time on the fringing reef flat was modeled to
28 decrease exponentially with increasing sea-level rise as the
29 magnitude of sea-level rise approached the mean water
30 depth over the reef flat. The model results presented here

suggest that a 0.5–1.0 m rise in sea level will likely 33
increase coastal erosion, mixing and circulation, the 34
amount of sediment resuspended, and the duration of high 35
turbidity on exposed reef flats, resulting in decreased light 36
availability for photosynthesis, increased sediment-induced 37
stress on the reef ecosystem, and potentially affecting a 38
number of other ecological processes. 39
40

Keywords Sea level · Fringing reef · Waves · Currents · 41
Sediment · Erosion 42

Introduction 43

While rising sea-surface temperatures and ocean acidifi- 44
cation have received most of the attention regarding the 45
impacts of climate change on coral reefs, the impact to coral 46
reefs from predicted future rising sea level has only 47
been addressed by a few researchers (Graus and Macintyre 48
1998; Ogston and Field 2010). A number of recent studies 49
(Grinsted et al. 2009; Merrifield et al. 2009) point out that 50
not only is global sea level rising, but the rate is increasing in 51
response to global climate change. Syntheses by Grinsted 52
et al. (2009) and Nicholls and Cazenave (2010) suggest that 53
global mean sea level in 2100 may exceed the 2000 level by 54
two to three times the average IPCC (2007) projection of 55
approximately 60 cm above 2000 levels. Since corals' 56
upward growth is constrained by exposure to air at low tides, 57
Buddemeier and Smith (1988) and Edwards (1995) sug- 58
gested that coral reef flats may benefit from the additional 59
accommodation space, as detrimental exposure to air would 60
decrease with sea-level rise. Buddemeier and Smith (1988), 61
however, qualify this conclusion as long as "... [coral] 62
communities are protected from destructive waves and not 63
subjected to heavy sedimentation...". 64

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65 Increased turbidity over coral reefs due to suspended
66 sediment can decrease light available for photosynthesis
67 (Marszalek 1981; Phillip and Fabricius 2003; Piniak and
68 Storlazzi 2008) and modify coral reef zonation by affecting
69 coral fertilization and recruitment, which, in turn, can
70 result in stress to or mortality of corals (e.g., Rogers 1990;
71 Phillip and Fabricius 2003; Fabricius 2005). Recently, a
72 number of studies have addressed the growing problem of
73 coral reefs impacted by anthropogenic modification of
74 coastal watersheds (e.g., Wolanski et al. 2003; Field et al.
75 2008), acknowledging that climate change may alter the
76 quantity and timing of sediment delivery to coral reefs.

77 Despite the widespread discussion of climate change
78 impacts on reefs, there has been little discussion of how
79 sea-level rise may affect fringing coral reefs, in terms of
80 both hydrodynamics and sediment dynamics. A number of
81 studies of bio-physical coupling on reefs (e.g., Sebens and
82 Johnson 1991; Edwards 1995; Falter et al. 2004; Storlazzi
83 et al. 2005) have shown that hydrodynamics control many
84 ecological aspects of reef systems, including photosyn-
85 thesis, nutrient uptake, prey capture, coral bleaching, and
86 species distribution. Ogston and Field (2010) presented
87 one-dimensional model results from Molokai, Hawaii,
88 USA, demonstrating that twenty-first-century sea-level rise
89 will increase wave heights and suspended-sediment con-
90 centrations and cause longer periods of elevated turbidity
91 on the coral reef flat.

92 In this paper, a two-dimensional numerical profile model
93 of hydrodynamics and sediment transport over the Molokai
94 fringing reef was calibrated with in situ data and was driven
95 by meteorologic and oceanographic forcing conditions that
96 characterize most exposed (not sheltered) tropical coral
97 reefs. The goal of this effort is to better understand the re-
98 lative importance of different processes (e.g., winds and
99 waves) to hydrodynamics and sediment transport, and the
100 contribution of these different characteristic sets of forcing
101 conditions to annual sediment fluxes. Model results for
102 various projections of sea-level rise are presented to identify
103 the relative importance of these different forcing conditions
104 to hydrodynamics and sediment transport in different sea-
105 level rise scenarios. Reef accretion or changes in roughness
106 were not modeled in these sea-level rise scenarios because
107 published vertical reef flat accretion rates for exposed
108 fringing reefs (1–4 mm year⁻¹ per Buddemeier and Smith
109 1988; Montaggioni 2005) are up to an order of magnitude
110 smaller than the rates of sea-level rise projected for the years
111 2000–2100 (8–16 mm year⁻¹ per Grinstead et al. 2009;
112 Nicholls and Cazenave 2010). These data suggest that pro-
113 jected sea-level rise will outstrip potential new reef flat
114 accretion, resulting in a net increase in water depth over
115 exposed fringing reef flats on the order of 0.4–1.5 m during
116 the twenty-first century. Lastly, the implications of these

results in the health and sustainability for fringing coral reefs
under projected sea-level rise are discussed.

Study area

The data presented here are from Molokai, Hawaii, in the
north-central Pacific Ocean (~21°N, 157°W) between the
islands of Oahu and Maui. The physical environment in
the Hawaiian Islands during the summer is dominated
by 5–10 m s⁻¹ northeasterly trade winds that generate
wave heights of 1–3 m with periods of 5–8 s and small
(1–2 m), long-period (14–25 s) south swells (Moberly and
Chamberlain 1964). Winter conditions, typically beginning
in October and extending through March, are character-
ized by storms and North Pacific swell that produce wave
heights of 3–6 m with periods of 10–18 s that approach
from the northwest. Due to shadowing by the surrounding
islands and the island of Molokai itself, however, most of
south Molokai's fringing reef is sheltered from large North
Pacific swell but is exposed to the other wave climates.
Hawaii has a mixed, semidiurnal microtidal regime, with
the mean daily tidal range of approximately 0.6 m and the
minimum and maximum daily tidal ranges are 0.4 and
0.9 m, respectively (Ogston et al. 2004; Storlazzi et al.
2004).

The morphology of the south Molokai fringing reef is
discussed in detail by Storlazzi et al. (2003) and is sum-
marized here. The reef flat, a roughly horizontal surface
with water depths ranging from 0.3 to 2.0 m, extends
seaward from the shoreline for distances from 0.5 to
1.5 km offshore. Calcareous marine sediment dominates
the coarse-grained fraction of the bed sediment (58–65%;
Field et al. 2008) across the entire fringing reef tract. The
inner portion of the reef flat is covered by a wedge of
muddy sand (80–90% of the silt and finer grain sizes are
terrigenous in origin) that pinches out roughly 200–300 m
offshore (Fig. 1a). From this point out to roughly 500 m
offshore, an ancestral reefal hardground is intermittently
exposed or mantled by sediment and algae. Shore-normal
ridge-and-runnel structure characterizes the reef flat from
500 m out to the reef crest roughly 1,000 m offshore. The
coral ridges are covered by low percentages of live coral,
and the runnel depressions are filled by calcareous sedi-
ment. The reef crest, where most deepwater waves break, is
generally well defined along many of the fringing reefs in
Hawaii and is locally covered by encrusting coralline algae
and robust lobate and encrusting corals. Offshore of the
reef crest, from depths of 3–30 m, lies the fore reef that is
generally characterized by 1–3 m high shore-normal spur-
and-groove structures covered by discontinuous, highly
variable percentages of live coral (Jokiel et al. 2001).

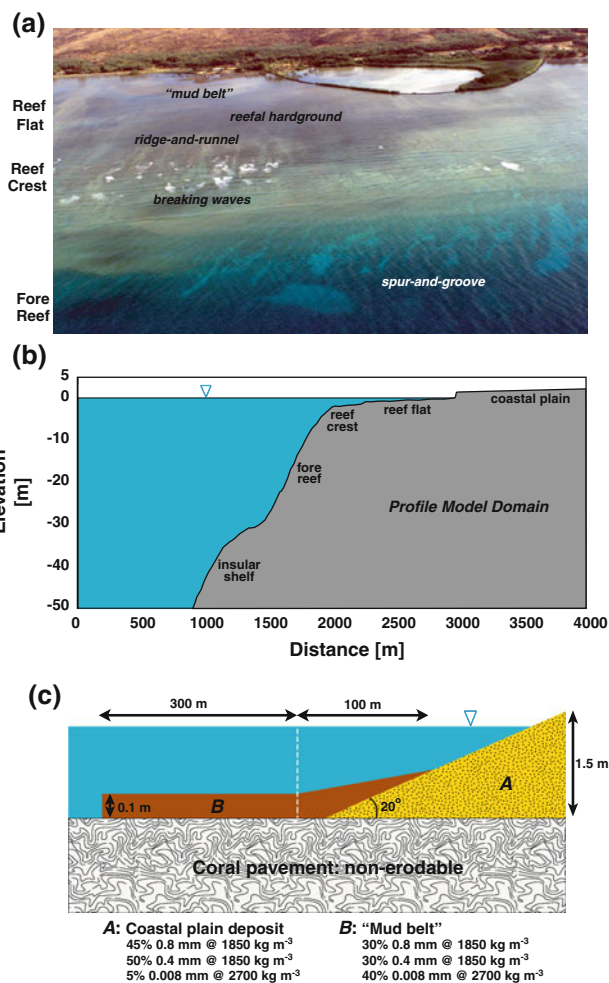


Fig. 1 The bathymetry, dominant zones, and geologic features of a fringing coral reef. **a** Oblique aerial photograph of the south Molokai, Hawaii, fringing coral reef. **b** Bathymetry and topography used in the numerical model. **c** The morphology and sedimentology of the inner reef flat, “mud belt”, and coastal plain deposit in the numerical model

166 **Methods**

167 **Field observations**

168 The in situ observations used for model calibration and
 169 validation were presented by Storlazzi et al. (2004) and
 170 Presto et al. (2006) and are summarized here. Current
 171 velocity data were collected via acoustic Doppler current
 172 profilers, acoustic Doppler velocimeters, or single-point
 173 electromagnetic current meters. Data loggers collected and
 174 stored data from these instruments as well as pressure to
 175 provide information on tides and waves. Optical back-
 176 scatter sensors and the acoustic backscatter from acoustic
 177 Doppler current profilers and acoustic Doppler velocime-
 178 ters provided information on turbidity and suspended-
 179 sediment concentrations. These instruments were mounted

on tripods or deployed from a mobile “backpack” in water 180
 depths ranging from 0.5 to 11.0 m. 181

Numerical modeling 182

The Delft3D Online Morphology system (Lesser 2004; 183
 Delft User Manual 2010) was used to model hydrody- 184
 namics and sediment transport over the south Molokai 185
 fringing coral reef. The main components are the two-way 186
 coupled Delft3D SWAN and FLOW modules modeling 187
 waves and currents, respectively. FLOW forms the core of 188
 the model system, simulating water motion due to tidal and 189
 meteorological forcing by solving the unsteady shallow- 190
 water equations that consist of the continuity equation, the 191
 horizontal momentum equations, and the transport equation 192
 under the shallow water and Boussinesq assumptions. 193
 Wave effects, such as enhanced bed shear stresses and 194
 radiation stresses, are included in the flow simulation by 195
 coupling the FLOW module with stationary runs of the 196
 third-generation SWAN wave model (Walstra et al. 2000). 197
 SWAN is based on discrete spectral action balance equa- 198
 tions, computing the evolution of random, short-crested 199
 waves (Holthuijsen et al. 1993; Booij et al. 1999; Ris et al. 200
 1999). Physical processes included are the generation of 201
 waves by wind, nonlinear quadruplet and triad wave-wave 202
 interactions, and dissipation due to whitecapping, bottom 203
 friction, and depth-induced breaking. 204

The Delft3D Online Morphology model was used to 205
 resolve the sediment resuspension and transport dynami- 206
 cally. At each computational time-step, Online Morphol- 207
 ogy supplements the FLOW module results with sediment 208
 transport using the Van Rijn (1993) formulation, wherein a 209
 distinction is made between bed-load and suspended-load 210
 transport. Bed-load transport represents the transport of 211
 sand particles in the wave boundary layer close to the 212
 seabed. Suspended-sediment transport is computed by the 213
 advection–diffusion solver. To describe sediment charac- 214
 teristics, additional formulations are included to account 215
 for density effects of sediment in suspension, settling 216
 velocity, vertical diffusion coefficient for sediment, sus- 217
 pended-sediment correction vector, and sediment exchange 218
 with the bed. The elevation of the bed is dynamically 219
 updated at each computational time-step by calculating the 220
 change in mass of the bottom sediment resulting from the 221
 gradients in sediment transport. 222

The bed was schematized in three non-cohesive sedi- 223
 ment classes (Fig. 1c) to represent the sediment observed 224
 on the reef flat off south Molokai (Field et al. 2008). For 225
 the fine terrigenous sediment (fine silt), a mean grain size 226
 (d_{50}) of 0.008 mm and a density of $2,700 \text{ kg m}^{-3}$ was 227
 prescribed, medium-sized carbonate sediment (fine sand) 228
 was characterized by $d_{50} = 0.2 \text{ mm}$ and density of 229
 $1,850 \text{ kg m}^{-3}$, and coarse carbonate sediment (medium 230

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231 sand) with $d_{50} = 0.4$ mm and density of $1,850 \text{ kg m}^{-3}$. In
 232 the model, sediment was initially present only in a narrow
 233 band extending 400 m from the shoreline while the
 234 remainder of the model was initially a non-erodible layer
 235 that represented the coral pavement of the fringing reef flat
 236 (thin, discontinuous sediment deposits were discounted)
 237 based on field observations. The total amount of sediment
 238 available in the profile was 56.7 m^3 of fine sediment,
 239 18.2 m^3 of which is contained in the mud belt (“B”),
 240 27.9 m^3 of medium sediment, and 784 m^3 of coarse sedi-
 241 ment contained in the coastal plain deposit (“A”). The
 242 volumes in the coastal plain deposit were somewhat arbi-
 243 trary, as these depended on the landward extent of the
 244 model. The model did allow for sediment accumulation
 245 and subsequent erosion of accumulated sediment on the
 246 coral pavement during the sediment-transport simulations.
 247 Sediment fractions are solved individually in the transport
 248 and bed-update modules and therefore were tracked sepa-
 249 rately. Hydraulic roughness length scales were varied
 250 between approximately 0.01 and 0.10 m, with the higher
 251 value used for coral surfaces based on previous observa-
 252 tions on Hawaiian reefs and numerical modeling results
 253 (e.g., Hearn 1999; Lowe et al. 2005), and the lower range
 254 (~ 0.01 m) set by the seabed grain size. Complete over-
 255 views of the formulations, testing, and validation of
 256 Delft3D Online Morphology have been reported in Lesser
 257 (2004). See Walstra and Van Rijn (2003), and Van Rijn
 258 (1993, 2007a, b, c) for the specific transport formulations.

259 Because most fringing reefs are relatively uniform
 260 alongshore in water depth and hydrodynamic roughness at
 261 larger scales (order ~ 100 s of m) but heterogeneous at
 262 smaller scales (order \sim m) due to ridge-and-runnel struc-
 263 tures on the reef flat and spur-and-groove structures on the
 264 fore reef, a fully realistic three-dimensional model at the
 265 spatial scales necessary to resolve the heterogeneity would
 266 be too computationally intensive. For this reason, a sche-
 267 matized two-dimensional profile model was used to acquire
 268 insight into the dominant sediment-transport processes
 269 across a fringing coral reef. Since the model is only one
 270 grid cell in the alongshore direction, high vertical and
 271 cross-shore resolution could be obtained while minimizing
 272 computational time. An underlying assumption in this
 273 approach is the dominance of wind- and wave-driven
 274 processes. This assumption is justified based on the small
 275 tidal velocities observed in the study area (Ogston et al.
 276 2004; Storlazzi et al. 2004; Presto et al. 2006). Model
 277 validation therefore was focused on accurate representation
 278 of the wave-breaking processes. Sensitivity analyses were
 279 performed on variations in the forcing by varying the
 280 seasonally schematized input conditions, profile dimen-
 281 sions, and the mean water level, including sea-level rise
 282 scenarios. The model bathymetry was based on averaging a
 283 number of cross-shore profiles from the high-resolution

284 SHOALS lidar data presented by Storlazzi et al. (2003).
 285 The model grid had a 10-m cross-shore resolution, result-
 286 ing in a total of 350 grid cells in the cross-shore direction
 287 for each vertical layer (Fig. 1b). The model was schema-
 288 tized in the vertical by 8 sigma layers with a thickness of 2,
 289 3, 5, 8, 12, 20, 25, and 25% of the total water depth from
 290 the seabed up to the surface.

291 High-resolution in situ measurements of tides, waves,
 292 currents, and suspended-sediment concentrations across the
 293 central part of the Molokai fringing reef were available
 294 only for a 40-d time frame (Storlazzi et al. 2004). These
 295 temporally limited in situ measurements were used for
 296 model calibration and validation; however, they may not be
 297 representative for the conditions that govern sediment
 298 transport over a range of seasons. To enable sediment-
 299 transport simulations for periods of time longer than the
 300 duration of instrument measurements, schematized forcing
 301 conditions for tides, wind, and waves were used. The
 302 technique described by Lesser et al. (2004) was applied to
 303 the water-level data from Presto et al. (2006) in order to
 304 generate a morphodynamic schematized tide to force the
 305 open ocean boundary (Table 1). The objective of the tidal
 306 schematization (input reduction) is to replace the full tide
 307 that is composed of all constituents that represent the full
 308 spring/neap cycle with a simplified 24.8-h tidal cycle that
 309 closely matches the residual tidally driven transport of the
 310 full lunar monthly tidal cycle. Such simplified tide should
 311 reproduce the residual sediment-transport rates and result-
 312 ing morphological change over the period of interest in the
 313 entire model domain.

314 The second important schematization was that of the
 315 wind and wave climate. The schematization used in this
 316 effort was based on the analysis of meteorologic and
 317 oceanographic data for the region (Presto et al. 2006;
 318 Storlazzi and Jaffe 2008) but is characteristic of most
 319 exposed coral reefs worldwide (e.g., Spalding et al. 2001;
 320 Riegl and Dodge 2008). This schematization resulted in
 321 four classes of distinct forcing conditions (Table 2). The
 322 Trade Wind conditions are the most prevalent, occurring
 323 62% of the time during a year ($226 \text{ days year}^{-1}$). During
 324 Trade Wind conditions, the wind is relatively strong and
 325 wave heights are moderate. Events characterized by minor
 326 wind and wave energy are schematized by the Variable

Table 1 Tidal constituents of the simplified tide used to force the open ocean model boundary based on the application of the Lesser et al. (2004) methodology

Constituent	Amplitude (m)	Phase (°)
AO	0.780	–
M2	0.178	64.7
C1	0.131	75.15

327 group, which occurs 25% of the time (91 days year⁻¹).
 328 High-energy Swell (large, long-period waves with weak
 329 winds) and Storm conditions (large waves and strong
 330 winds) occur less frequently (10% [37 days year⁻¹] and
 331 3% [11 days year⁻¹], respectively). The model was then
 332 run for a year's time driven by the four sets of wind and
 333 wave forcing conditions for their respective durations
 334 (Trade Wind for 226 days, Storm for 11 days, etc.).

335 Results

336 Model calibration and validation

337 Calibration and validation of the numerical model focused
 338 on the water levels and wave heights. While the small tidal
 339 velocities were of minor importance in these simulations,
 340 the tidal water levels needed to be represented accurately.
 341 Root-mean-squared (RMS) errors between measured and
 342 modeled water levels across the reef (Fig. 2a–c) were
 343 0.014 m, and the RMS errors in wave height across the reef
 344 were 0.017 m (Fig. 2d–f), showing very good correspon-
 345 dence between the model and in situ data (Storlazzi et al.
 346 2004). The modeled mean current speeds (7.8 cm s⁻¹) on
 347 the reef flat slightly exceeded the observed values of
 348 5.3 ± 3.7 cm s⁻¹ (mean \pm SD; not shown); the modeled
 349 mean current speed on the fore reef (5.0 cm s⁻¹) corre-
 350 sponded well to the observations (4.1 ± 4.9 cm s⁻¹).
 351 Importantly, the modeled tidal and total (wind +
 352 wave + tide) current speeds were on the same order (0–5
 353 and 5–20 cm s⁻¹, respectively) as the measured currents
 354 and also showed the same proportion of greater (\sim 2–8
 355 times) alongshore current speeds on the reef flat and over the
 356 fore reef compared to the cross-shore current speeds. This is
 357 in contrast, however, to most observations and models of
 358 primarily cross-shore flow and sediment transport on atolls
 359 and barrier reefs (e.g., Hearn 1999; Hearn and Atkinson
 360 2000; Lowe et al. 2005) where vigorous wave-driven
 361 onshore flow over the reef flat can occur because it is bal-
 362 anced by strong return flow out of channels in the reef. In
 363 fringing reefs without a nearshore gully, wave-driven setup
 364 along the shoreline offsets this cross-reef flow and results in
 365 primarily alongshore flow and transport.

Contributions of waves and currents to sediment 366
 transport 367

368 With confidence that the numerical model was successfully 368
 369 reproducing the hydrodynamics on the reef flat and fore 369
 370 reef, the four schematized forcing conditions and the 370
 371 resulting sediment dynamics were modeled. The goal of 371
 372 these sediment dynamics simulations was not to reproduce 372
 373 reality, for the schematized model domain and duration of 373
 374 forcing conditions when compared to the limited in situ 374
 375 measurements made this not possible. Rather, the goal was 375
 376 to understand the relative contribution of the different 376
 377 forcing mechanisms to flow and sediment transport and 377
 378 how these contributions would vary with sea-level rise. The 378
 379 mean forcing and resulting suspended-sediment concen- 379
 380 trations for the four schematized forcing conditions are 380
 381 shown in Fig. 3; the resulting sediment transport for a 381
 382 1-year simulation comprised of the four schematized 382
 383 forcing conditions is shown in Fig. 4. Overall, sediment 383
 384 transport was dominated by the fine-grained fractions. No 384
 385 coarse- or medium-grained fractions were moved on the 385
 386 reef flat during the base (sea level = 0.00 m) simulations. 386
 387 The bulk of the wave energy was dissipated along the reef 387
 388 crest, and the depth-averaged current speeds in the “mud 388
 389 belt” were small and did not exceed the critical threshold 389
 390 of motion for the larger (medium and coarse sand) grain- 390
 391 size fractions, resulting in a narrow band of elevated fine- 391
 392 grained sediment concentrations along the shoreline. No 392
 393 significant sediment losses to deep water were encountered 393
 394 as a result of the minor water level gradient-induced off- 394
 395 shore flow near the bed, resulting in most of the sediment 395
 396 transport laterally alongshore in a band extending from the 396
 397 shoreline seaward approximately 400 m, which matches 397
 398 observations (Presto et al. 2006; Field et al. 2008). 398

399 The model results show distinctively different transport 399
 400 rates during the different forcing conditions. Storm con- 400
 401 ditions dominated the hydrodynamics and resulting sedi- 401
 402 ment dynamics despite their low frequency of occurrence. 402
 403 Storm conditions generated the greatest setup on the reef 403
 404 flat (Fig. 3b) due to the strong winds and highest wave 404
 405 energy on the fore reef (Fig. 3c). The winds and waves 405
 406 drove strong (\sim 5–15 cm s⁻¹) currents across the fore reef 406
 407 and reef flat (Fig. 3d) and, together, generated high shear 407

Table 2 Model schematization of wind and wave conditions. In the model, the coast trends north–south (0–180°), with eastward (90°) being oriented onshore

Climate	Percent of days year ⁻¹	Wind direction (°)	Wind speed (m s ⁻¹)	Wave height (m)	Wave period (s)	Wave direction (°)
Trade Wind	62	190	10	1	6	190
Variable	25	80	3	0.5	6	240
Swell	10	170	3	1	14	280
Storm	3	280	20	1.5	8	280

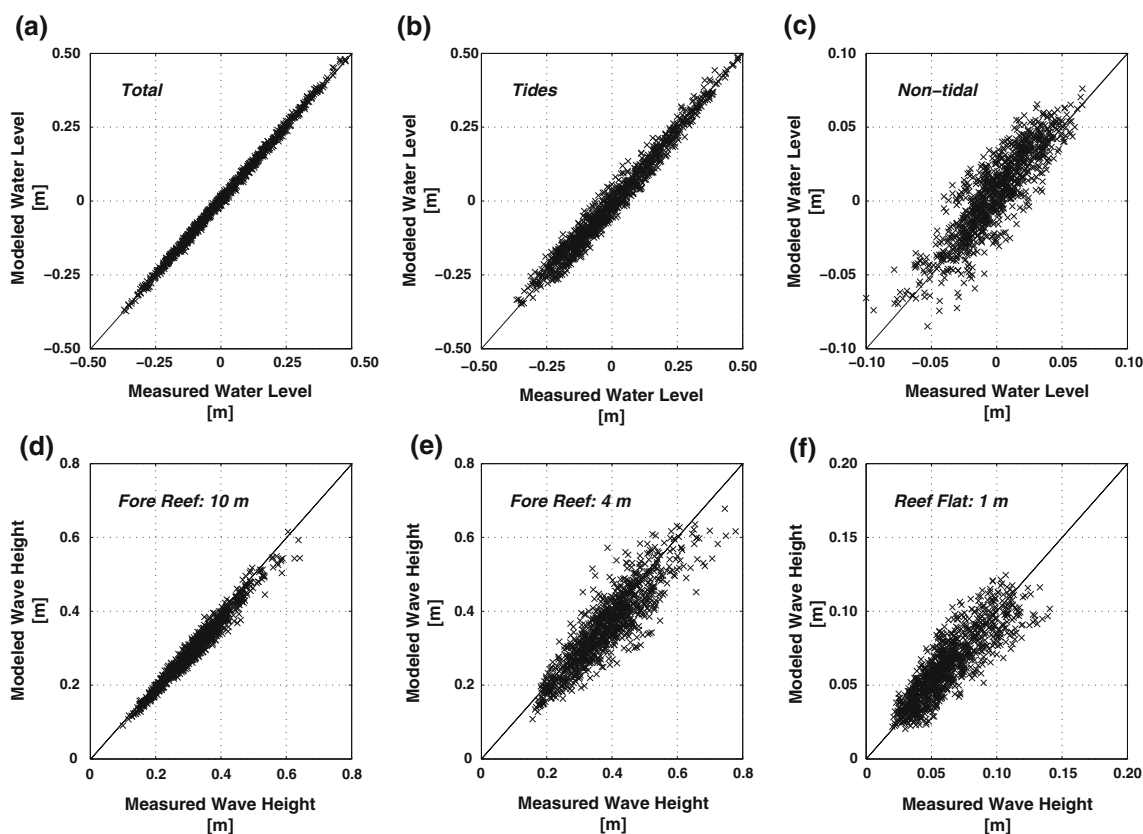


Fig. 2 Comparison of in situ measurements (Storlazzi et al. 2004) and model results of water levels and wave heights. **a** Total water level on the reef flat (depth ~ 1 m). **b** The tidal component of water level on the reef flat. **c** The non-tidal component of water level on the

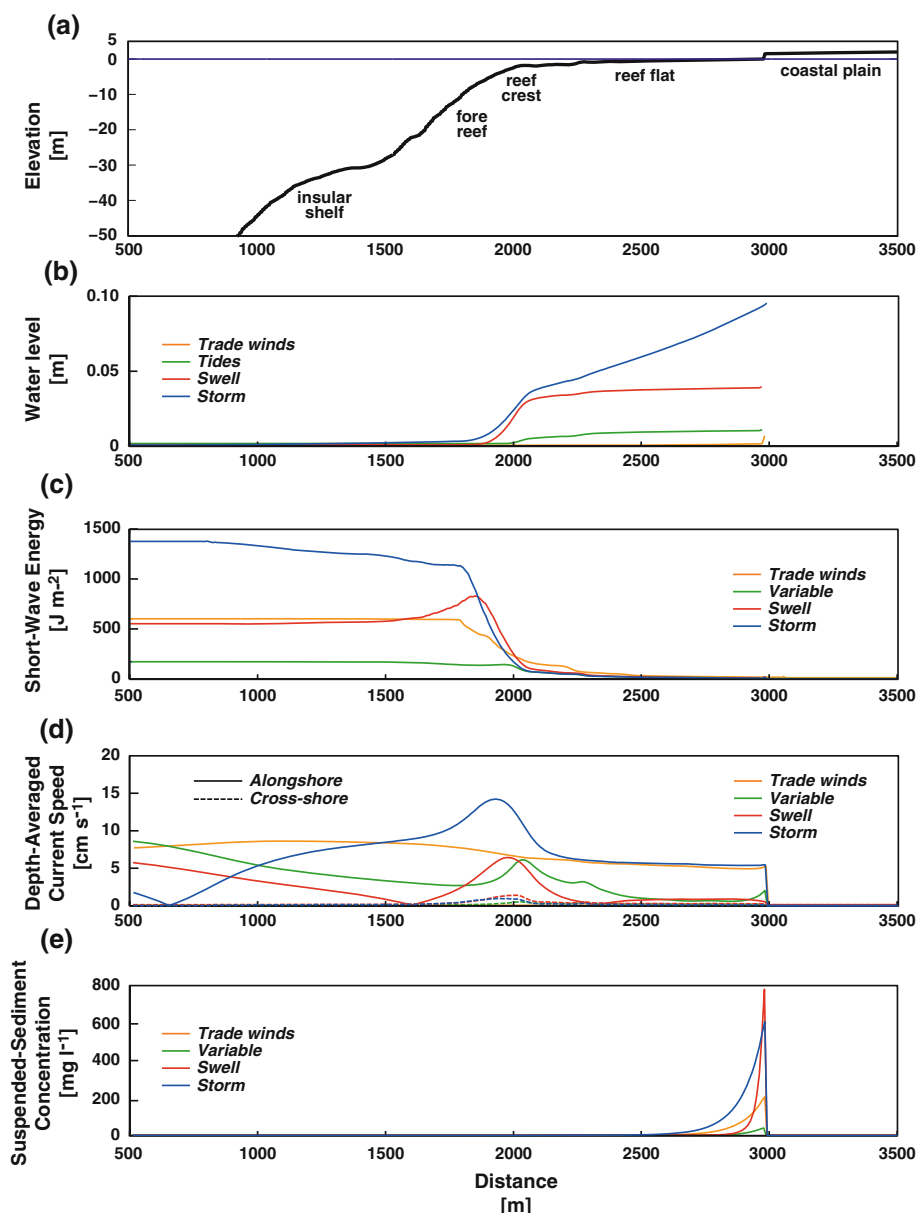
reef flat. **d** Wave height on the fore reef (depth ~ 10 m). **e** Wave height on the fore reef (depth ~ 4 m). **f** Wave height on the reef flat (depth ~ 1 m). These comparisons show that the errors between observed and modeled water levels and wave heights are less than 5%

408 stresses that resuspended large quantities of sediments
 409 across the inner reef flat (Fig. 3e). Even though the Swell
 410 conditions did not generate strong currents, the long period
 411 of the waves generated substantial long-wave energy that
 412 caused almost 5 cm of setup. These long-wave motions
 413 during the Swell conditions resulted in higher peak sus-
 414 pended-sediment concentrations right at the shoreline than
 415 modeled during the Storm conditions, but this zone of
 416 elevated suspended-sediment concentrations was confined
 417 to closer to shore and thus a lower total mass of sediment
 418 in suspension over the reef flat than in the Storm conditions.
 419 The Trade Wind conditions generated relatively strong
 420 wind-driven currents with little wave forcing and resulted
 421 in elevated suspended-sediment concentrations close to the
 422 shoreline, but these concentrations were on the order of a
 423 third to a quarter of those modeled during Storm and Swell
 424 conditions. Lastly, Variable conditions generated relative
 425 weak currents across the reef flat that resulted in sus-
 426 pended-sediment concentrations on the order of a third to a
 427 quarter of those modeled during Trade Wind conditions
 428 and almost an order of magnitude lower than during Storm
 429 and Swell conditions. Importantly, the waves, current

430 speeds, and resulting suspended-sediment concentrations
 431 during the modeled forcing conditions match well with the
 432 in situ data collected under similar atmospheric and
 433 oceanographic forcing (Ogston et al. 2004; Storlazzi et al.
 434 2004; Presto et al. 2006).

435 The relative contribution of the different sets of forcing
 436 conditions to annual sediment flux (Fig. 4) shows that
 437 Storm conditions are the dominant contributor to annual
 438 sediment flux, contributing just over twice the sediment
 439 flux that was modeled during Trade Wind conditions but in
 440 only 5% of the time (Fig. 4c; Table 3). This high per-
 441 centage (63%) of total annual sediment flux in only 11 d
 442 shows the importance of not only large wave and strong
 443 winds generating strong currents and high shear stresses,
 444 but the importance of setup increasing water depth over the
 445 reef flat that, in turn, allows for larger waves and stronger
 446 currents by reducing the hydrodynamic roughness relative
 447 to the depth of the water column. While Swell conditions
 448 resulted in high suspended-sediment concentrations at the
 449 shoreline (Fig. 3e), the greater cross-shore extent of ele-
 450 vated turbidity during Trade Wind conditions and their
 451 more frequent occurrence resulted in just under an order of

Fig. 3 Modeled cross-shore variations in hydrodynamics and sediment dynamics for the four forcing conditions. **a** Morphology of the model domain. **b** Water level. **c** Short-wave energy. **d** Depth-averaged current speeds, with alongshore currents as *solid lines* and cross-shore currents as *dashed lines*. **e** Suspended-sediment concentrations. While Swell conditions result in the highest wave energy and suspended-sediment concentrations on the reef flat, the greatest volume of suspended sediment over the reef flat results from Storm conditions



452 magnitude greater contribution to the annual sediment flux
 453 than Swell conditions (Fig. 4c; Table 3). Also of note is the
 454 relative contributions of suspended and bed-load flux to the
 455 annual total sediment flux (Fig. 4c). While the total sedi-
 456 ment flux during Trade Wind conditions is primarily
 457 material in suspension, the more energetic Storm condi-
 458 tions result in greater erosion of the “mud belt” deposit and
 459 a resulting greater proportion of bed load to the annual total
 460 sediment flux.

461 Effects of sea-level rise on waves and currents

462 With confidence that the numerical model was successfully
 463 reproducing the hydrodynamics and sediment dynamics on
 464 the reef flat under present conditions (sea level = 0.00 m),

mean sea level was then elevated (+0.10, +0.25, +0.50, 465
 and +1.00 m) to investigate the influence of sea-level rise 466
 on waves across the fringing reef. Since the four sets of 467
 forcing conditions and 5 different sea-level rise scenarios 468
 result in 20 different model runs, for visualization purposes 469
 the annual weighted (by frequency of occurrence during 470
 the year) mean hydrodynamics and resulting sediment 471
 dynamics for the four different forcing conditions are 472
 presented for the 5 sea-level rise scenarios in Fig. 5. As sea 473
 level increased, the breaking wave height at the reef crest 474
 decreased and the location of maximum wave breaking (as 475
 evidenced from wave energy dissipation) moved landward 476
 (Fig. 5b, c) as more wave energy was able to propagate up 477
 onto the reef flat, resulting in greater wave heights on the 478
 reef flat. The depth-limited nature of wave height on the 479

Fig. 4 Schematization of the forcing data and the resulting modeled net annual sediment flux. **a** Variations in wave heights for the forcing conditions. **b** Variations in wave periods for the forcing conditions. **c** Variations in wind speeds for the forcing conditions. **d** Total net suspended-sediment flux. Storm conditions are the dominant contributor to annual sediment flux, contributing just over twice the sediment flux that was modeled during Trade Wind conditions but in only 5% of the time

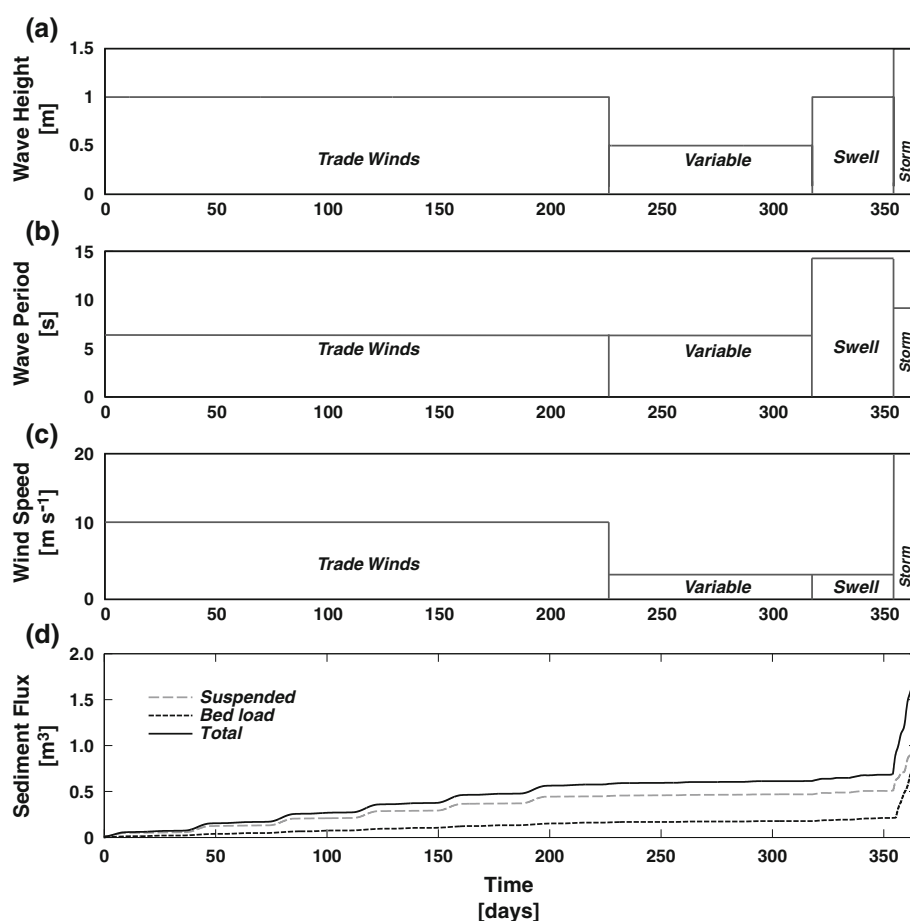


Table 3 Modeled cumulative total annual sediment transport and residence times under present and predicted future sea-level rise scenarios

Sea level (m)	Cumulative transport (m ³)					Residence time (years)
	Trade Wind	Variable	Swell	Storm	Total	
0.00 (present)	0.570	0.031	0.073	1.146	1.820	10.1
+0.10	1.037	0.056	0.105	1.315	2.513	7.3
+0.25	1.942	0.111	0.141	1.454	3.647	5.0
+0.50	3.608	0.207	0.132	2.751	6.697	2.7
+1.00	1.047	0.311	0.749	24.478	26.585	0.7

480 reef flat is evident not only in the landward decrease in
481 wave height and energy dissipation due to wave breaking,
482 but also how both of these parameters increase with
483 increasing sea level. As sea level increased, the propaga-
484 tion of larger waves over the reef crest onto the reef flat and
485 in situ growth of wind-waves on the reef flat resulted in
486 elevated peak-bed wave-induced shear stresses (Fig. 5d),
487 especially close to shore in the shallows where the “mud
488 belt” contains a significant proportion of terrestrial materi-
489 al. While increased sea level resulted in increased wave

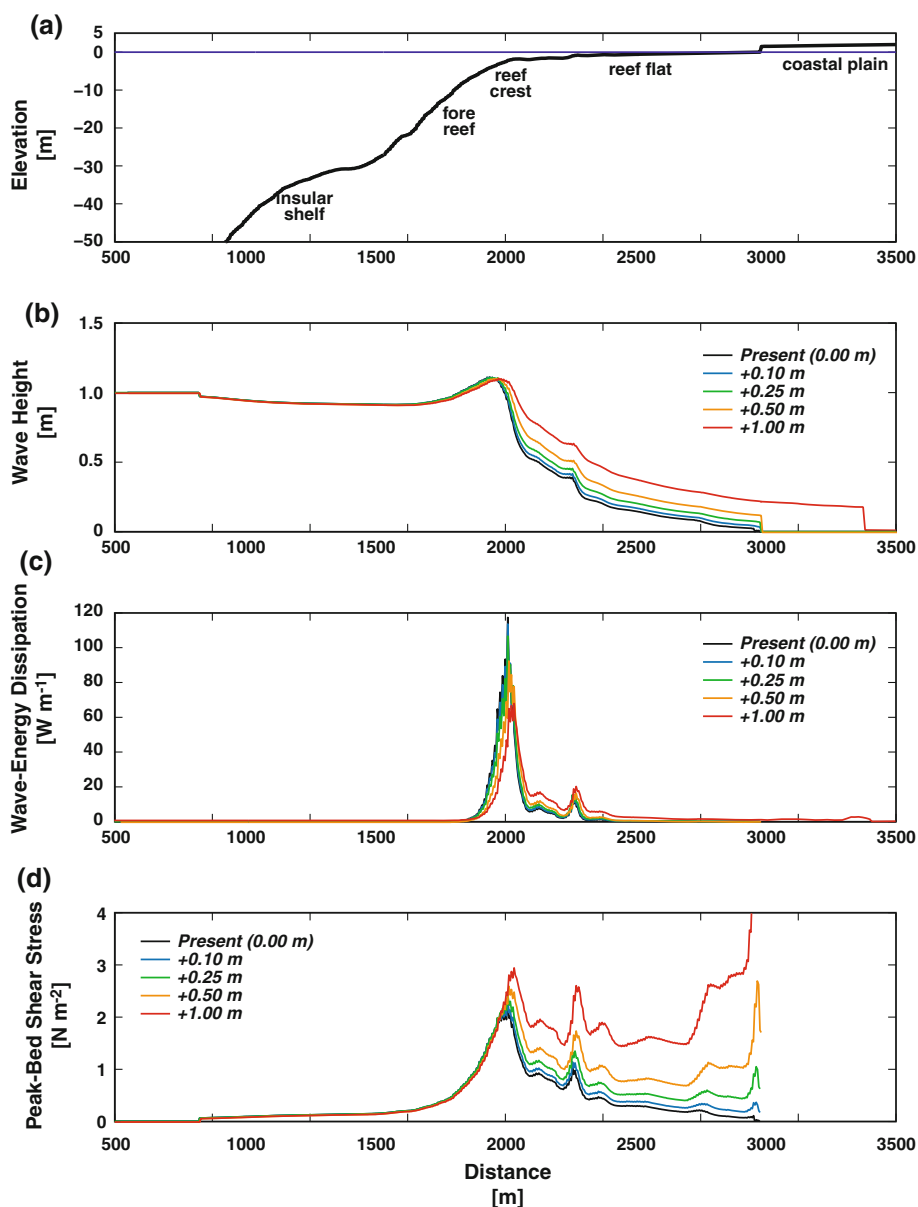
heights, dissipation, and peak-bed shear stresses on the reef 490
flat, the maximum radiation stress-induced setup on the 491
reef flat due to wave breaking decreased and elevated setup 492
extended farther offshore toward the reef crest (Fig. 6b), 493
possibly due to an increase in flow depth over the reef crest 494
and reef flat relative to the hydrodynamic roughness 495
imparted by the corals and ridge-and-runnel structure. At a 496
sea level +1.00 m (almost doubling water depth over much 497
of the reef flat), however, run-up onto the coastal plain 498
extended to +0.30 m above the oceanic water level 499
(+1.30 m total). 500

Effects of sea-level rise on sediment dynamics 501

Similar to the study of hydrodynamics, the effect of sea- 502
level rise on sediment-transport rates was investigated by 503
elevating mean sea level (+0.10 m, +0.25 m, +0.50 m, 504
and +1.00 m). The remainder of the model schematiza- 505
tions and parameter settings were unchanged compared to 506
the base case simulation (sea level = 0.00 m). Sediment 507
transport in the sea-level rise scenario model runs was 508
governed by the fine sediment fractions, similar to the 509
present-day (sea level = 0.00 m) model runs. Current 510
speeds in the “mud belt” were relatively small 511

Fig. 5 Modeled cross-shore variations in wave parameters as a function of water depth.

a Morphology of the model domain. **b** Wave height. **c** Wave energy dissipation. **d** Wave-induced peak-bed shear stress. Wave height, energy dissipation, and peak-bed shear stress decrease at the reef crest but increase in the reef flat with increasing water depth. Note the greater cross-shore extent (distance $\sim 3,000$ – $3,350$ m) of the parameters at a sea-level rise of 1.0 m due to approximately 350 m of erosion into the coastal plain deposit by the larger waves and resulting high shear stresses

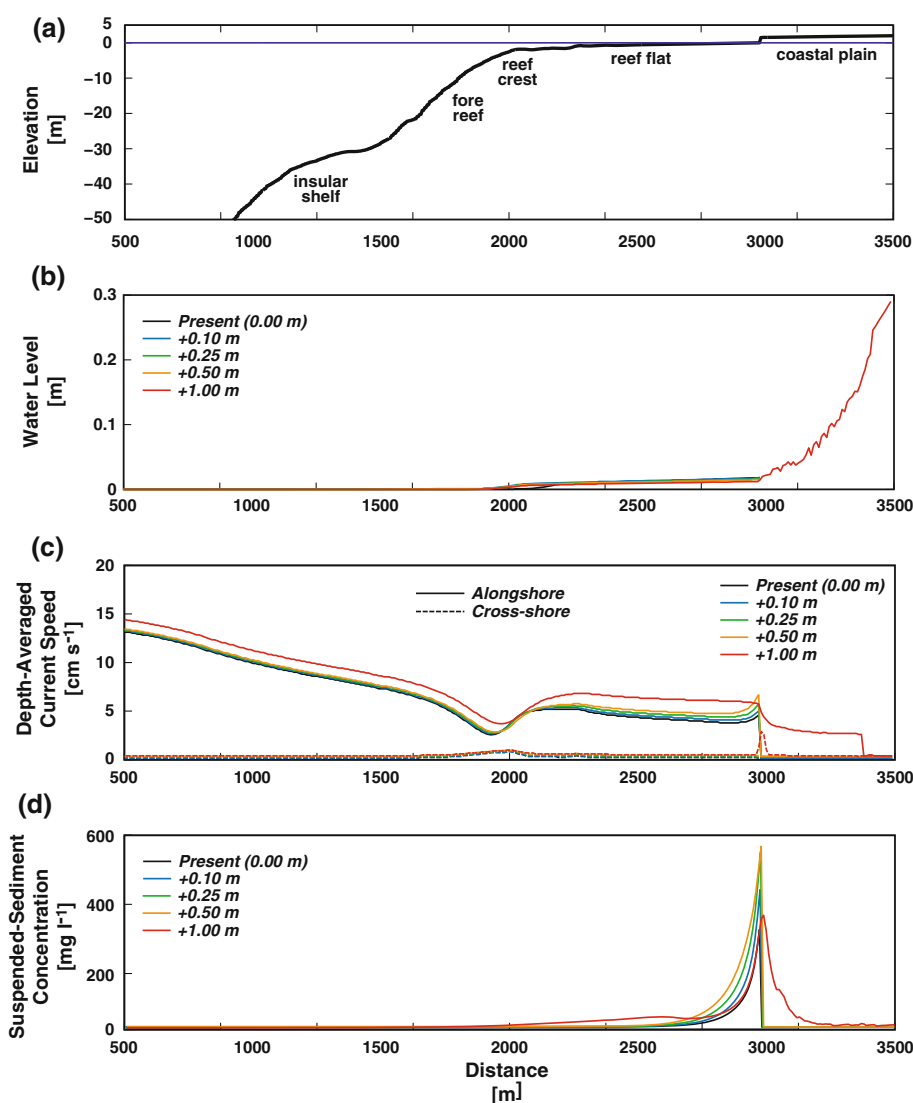


512 ($\sim 5 \text{ cm s}^{-1}$; Fig. 6c) and did not exceed the critical
513 threshold of motion for the larger (sand-sized) grain-size
514 fractions. As sea level was increased, current speeds
515 increased and the relative minima in alongshore current
516 speeds and maxima in cross-shore current speeds at the reef
517 crest due to wave breaking migrated onshore, similar to the
518 maxima in wave height and energy dissipation (Fig. 5).

519 The higher wave-induced high peak-bed shear stresses
520 (Fig. 5d) and current speeds (Fig. 6c) close to shore
521 resulted in a narrow band of suspended-sediment concen-
522 trations and transport along the shoreline (Fig. 6d).
523 Transport rates due to Storm conditions were an order of
524 magnitude larger than the Swell conditions, which were an
525 order higher than Trade Wind conditions. While suspen-
526 ded-sediment concentrations increased in magnitude

527 close to the shoreline and elevated suspended-sediment
528 concentrations extend further across the inner half of the
529 reef flat with increasing sea level between +0.10 m and
530 +0.50 m, there is a distinct change in this pattern when sea
531 level reached +1.00 m. When sea level was set at +1.00 m
532 (almost doubling water depth over much of the reef flat),
533 enough deep-water wave energy was able to propagate
534 onto the reef flat such that larger waves impacted the
535 shoreline, causing significant (>0.30 m) setup along the
536 shoreline. These waves eroded approximately 350 m hori-
537 zontally into the coastal plain deposit, resulting in wave
538 heights (Fig. 5b), water levels (Fig. 6b), currents (Fig. 6c),
539 and suspended-sediment concentrations (Fig. 6d) shore-
540 ward of the original shoreline location in the model at a
541 cross-shore distance of 3,000 m. The erosion of the coastal

Fig. 6 Modeled cross-shore variations in water level, current speeds, and suspended-sediment concentrations as a function of water depth. **a** Morphology of the model domain. **b** Total water level, including wave-induced setup. **c** Alongshore (solid line) and cross-shore (dashed line) current speeds. **d** Suspended-sediment concentrations. Wave height, energy dissipation, and peak-bed shear stress decrease at the reef crest but increase in the reef flat with increasing water depth. Note the greater cross-shore extent (distance $\sim 3,000$ – $3,350$ m) of the parameters at a sea-level rise of 1.0 m due to approximately 350 m of erosion into the coastal plain deposit by the larger waves (Fig. 5)

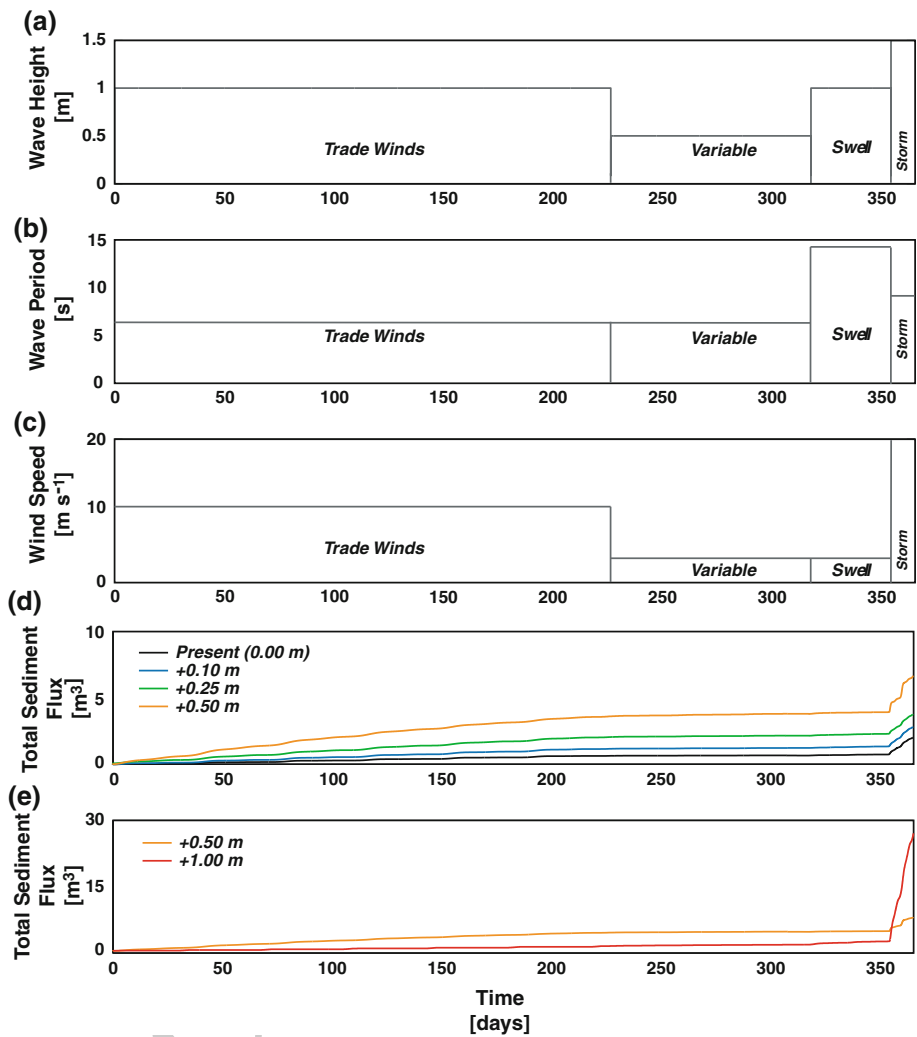


542 plain deposit resulted in a lower maximum suspended-
 543 sediment concentrations but a greater overall volume of
 544 material in suspension over the profile (Table 3) as the
 545 elevated suspended-sediment concentrations extended
 546 seaward out to the reef crest (cross-shore distance of
 547 2,000 m) and shoreward to the new shoreline (cross-shore
 548 distance of $\sim 3,350$ m). The lower maximum suspended-
 549 sediment concentrations at +1.00 m of sea-level rise
 550 resulted from the erosion of the primarily coarse-grained
 551 coastal plain deposit, which provided only 5% fine-grained
 552 material by volume to the reef flat that could easily be
 553 resuspended by the waves and currents.

554 Total sediment-transport rates for each of the simulations
 555 are displayed in Fig. 7d-e. There was a statistically
 556 significant exponential increase in sediment-transport rates
 557 with increasing sea-level rise ($r^2 = 0.999$ for $n = 5$;
 558 $P < 0.001$). As the water level rises over the reef flat, more
 559 wave energy propagated to the shoreline, resulting in a

560 greater energy transfer, higher bed shear stresses, greater
 561 sediment resuspension, and higher sediment-transport
 562 rates. Higher sea-level rise (0.50–1.00 m), which is on the
 563 order of the water depth over the reef flat, resulted in an
 564 increase in the dominance of Storm-type conditions due to
 565 larger waves breaking closer to the shoreline (Table 3). An
 566 estimate of the sediment residence time can be obtained by
 567 analyzing the sediment fluxes and available volume
 568 (Table 3). For this, the “mud belt” material (fine silt and
 569 fine sand) on the reef flat (Fig. 1a, c) was the primary
 570 focus, as medium sand was generally not transported in the
 571 simulations. Annual cross-shore losses ($0.189 \text{ m}^3 \text{ year}^{-1}$)
 572 in the model were small (10.4%) compared to the along-
 573 shore transport ($1.820 \text{ m}^3 \text{ year}^{-1}$). The residence time of
 574 sediment in the model can be obtained by analysis of the
 575 cumulative transport. At present (0.00 m) based on a loss
 576 of $1.820 \text{ m}^3 \text{ year}^{-1}$, it would take 10.1 years for the fine
 577 silt-sized terrestrial sediment to be completely removed

Fig. 7 Schematization of the forcing data and the resulting modeled net annual sediment flux for the five sea-level rise scenarios. **a** Variations in wave heights for the forcing conditions. **b** Variations in wave periods for the forcing conditions. **c** Variations in wind speeds for the forcing conditions. **d** Total net suspended-sediment flux for the lower magnitudes of sea-level rise. **e** Total net suspended-sediment flux for the higher magnitudes of sea-level rise. Note that subplots “d” and “e” show some of the same data, but have different y-axes to highlight the details in the time series. Higher sea-level rise (0.50–1.00 m) on the order of the water depth over the reef flat resulted in an increase in the dominance of Storm-type conditions to total sediment flux due to larger waves breaking closer to the shoreline



578 from the “mud belt” in the profile without any additional
 579 inputs. Sediment residence time on the reef flat displayed a
 580 statistically significant exponential decrease with increas-
 581 ing sea-level rise ($r^2 = 0.999$ for $n = 5$; $P < 0.001$) as the
 582 waves and current speeds increase. These calculations are
 583 assumed to be a lower limit of the sediment residence time;
 584 Presto et al. (2006) predicted a residence time of approx-
 585 imately 30 years based on in situ measurements. For the
 586 sand-sized carbonate sediment, the sediment residence time
 587 varies around 644 years (rates $\sim 0.044 \text{ m}^3 \text{ year}^{-1}$) due to
 588 the very low frequency of mobilization.

589 **Discussion**

590 The numerical modeling of hydrodynamics and sediment
 591 transport over fringing coral reefs presented here suggest
 592 the following changes are expected to occur under future
 593 sea-level rise scenarios:

Waves

594
 595 Greater water depths over a fringing reef would reduce
 596 bottom friction and increase water depth relative to the
 597 wave height, resulting in larger and more energetic waves
 598 that could propagate over the reef crest and reef flat without
 599 breaking and larger wind-waves develop in situ on the reef
 600 flat, similar to the model results presented by Hearn (1999)
 601 and Hearn and Atkinson (2000). These findings are sup-
 602 ported by data from Storlazzi et al. (2004), who showed
 603 that while wave heights offshore of the reef crest on the
 604 fore reef (depth $\sim 10 \text{ m}$) are independent of sea level
 605 ($r^2 = 0.003$, $n = 961$, P not significant), both wave height
 606 and wave period on the reef flat (depth $\sim 1 \text{ m}$) are sig-
 607 nificantly correlated with sea level ($r^2 = 0.791$ and 0.797 ,
 608 respectively; both $n = 961$ and $P < 0.001$), suggesting that
 609 waves on the reef flat are depth-limited. As sea-level rise
 610 increases, the larger waves over the reef and the landward
 611 migration of the zone of primary incident wave breaking

Author Proof

612 will also modify the zone of high turbulence, primarily
613 moving it shorewards.

614 Currents

615 Increased water depth would result in stronger currents all
616 across the reef due to greater wave-driven flows from lar-
617 ger waves and the reduced height of hydrodynamic
618 roughness relative to water depth that would allow for
619 faster wind-driven currents to develop. The greatest
620 increases in current velocity would be in shallow water on
621 the inner reef flat where the water depth is on order of the
622 hydrodynamic roughness of the seabed. This finding is
623 supported by data from Presto et al. (2006), who showed
624 the current speed at a given location on the reef flat was
625 statistically greater (mean difference = 2 cm s^{-2} , $n = 27$,
626 $P < 0.05$) during periods with higher sea level than during
627 periods of lower sea level. This would result in greater
628 water exchange across and lower residence time of water
629 on the reef flat, potentially altering the physical and
630 chemical properties of the water column. The increased
631 mixing and flushing of the reef flat with sea-level rise may
632 help to dilute material delivered to the inner reef flat from
633 the adjacent land, but it might also result in greater trans-
634 port of terrestrial sediment onto the fore reef. The model
635 results presented here, along with the observations made by
636 Ogston et al. (2004), Storlazzi and Jaffe (2008), and Lowe
637 et al. (2009) and modeled by Gourlay (1996), Hearn
638 (1999), and Hearn and Atkinson (2000), show the effect of
639 sea level on the magnitude of currents, driven both by wind
640 and by wave-breaking, on coral reef flats.

641 Sediment dynamics

642 An increase in wave energy and circulation due to elevated
643 sea level will also affect sediment dynamics across a
644 fringing coral reef. Larger waves resulting from high water
645 levels will generate increased wave-induced stresses,
646 which, in turn, will result in greater resuspension of sedi-
647 ment across the reef for a given grain size or composition
648 (e.g., density) as suggested by Ogston and Field (2010).
649 Statistically greater suspended-sediment concentrations
650 (mean difference = 46 mg l^{-1} , $n = 32$, $P < 0.001$) were
651 observed by Presto et al. (2006) during periods with higher
652 sea level; Storlazzi et al. (2004; Fig. 9 therein) showed that
653 suspended-sediment concentrations on both reef flat
654 ($r^2 = 0.383$ for $n = 961$; $P < 0.001$) and fore reef
655 ($r^2 = 0.238$ for $n = 961$; $P < 0.02$) were significantly
656 greater during periods with higher sea level. These obser-
657 vations, combined with their observations showing that
658 greater offshore flow occurred with higher sea level,
659 resulted in a statistically significant greater flux of sediment
660 off the reef flat ($r^2 = 0.369$ for $n = 961$; $P < 0.05$) and

661 over the fore reef ($r^2 = 0.576$ for $n = 961$; $P < 0.001$)
662 with higher sea levels. The greater resuspension and
663 transport would result in higher and longer persistence of
664 turbidity as the increased shear stresses and turbulence
665 would inhibit sediment from settling. The increased
666 resuspension and larger wave-orbital velocities with sea-
667 level rise may also alter patterns of abrasion of corals
668 adjacent to sedimentary deposits such as the sediment-filled
669 “grooves” of spur-and-groove structures. Although
670 alongshore current speeds on the reef flat and over the fore
671 reef are generally much greater (~ 2 – 8 times) than the
672 cross-shore current speeds (Ogston et al. 2004; Presto et al.
673 2006), there is strong coupling between offshore flow and
674 high suspended-sediment concentrations on reef flats such
675 that the greatest sediment fluxes generally have an offshore
676 component (Storlazzi et al. 2004). Storlazzi and Jaffe
677 (2008) showed similar reef flat–fore reef coupling off west
678 Maui, especially during periods of large waves and storms
679 when water levels are elevated over the reef flat due to
680 wind- and wave-induced setup.

681 As pointed out by Graus and Macintyre (1998) and
682 Ogston and Field (2010), the larger waves on the reef flat
683 that would result from sea-level rise will also increase the
684 delivery of energy to the coastline. As these larger, more
685 energetic waves reach the shoreline, which at present is in
686 quasi-equilibrium with the current wave climate, they
687 would exceed the critical shear stresses for resuspension of
688 the beach and coastal plain material, causing coastal ero-
689 sion and adding additional sedimentary material to the reef
690 flat, similar to the observations by Sheppard et al. (2005).
691 This additional material, resuspended by larger waves and
692 stronger currents, would likely exacerbate turbidity not
693 only on the reef flat but also likely on the fore reef as well
694 (Storlazzi et al. 2004; Storlazzi and Jaffe 2008). Although
695 the stronger currents may reduce the overall residence time
696 of any given sedimentary particle on the reef flat, the
697 increased supply of material by erosion and the increased
698 duration of resuspension for a given set of waves and
699 currents could potentially result in greater exposure of
700 corals to suspended sediment on both reef flat and fore reef.

701 The one-dimensional modeling by Ogston and Field
702 (2010) on waves and sediment resuspension and the two-
703 dimensional coupled hydrodynamics and sediment-trans-
704 port modeling presented here provide insight into the
705 potential affects of sea-level rise on flow and sediment
706 dynamics over an exposed fringing coral reef based on
707 current observations of forcing conditions (winds and
708 waves). The hydrodynamic and sediment-transport data
709 presented here suggest that while some protected fringing
710 coral reefs many benefit from the additional accommoda-
711 tion space as suggested by Edwards (1995), all will
712 undergo a number of changes in both chemical and bio-
713 logical processes (e.g., Sebens and Johnson 1991; Edwards

714 1995; Falter et al. 2004; Storlazzi et al. 2005) due to
715 changes in the hydrodynamics caused by sea-level rise.
716 Many exposed fringing coral reefs may be threatened by
717 additional input and resuspension of terrestrial sediment
718 that will likely negatively affect corals and the ecosystems
719 they support on both reef flat and fore reef, as first postu-
720 lated by Buddemeier and Smith (1988).

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